

# Intermolecular Electron Transfer Reactivity Determined from Cross-Rate Studies

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## ABSTRACT

Electron-transfer cross-reactions between neutral molecules and their radical cations spanning a wide range of structural type and intrinsic reactivity have been analyzed using classical Marcus theory. The principal factor found to govern intrinsic reactivity is the inner-shell bond reorganization energy. The HOMO–LUMO overlap of alkyl groups on reacting molecules is generally sufficient to provide facile electron transfer; however, a significant steric effect on this overlap is observed for hydrazines with alkyl groups larger than methyl.

## Introduction

Relationships between structure and reactivity form the core of our understanding of chemistry and are fundamental in allowing prediction of which compounds to employ for various functions. A basic determining factor of reactivity is the driving force for a reaction,  $\Delta G^0$ . Most structural changes alter both the driving force and intrinsic reactivity—the reactivity at constant driving force. Especially well-studied series that cover very large ranges in reactivity include solvolyses in polar solvents and acid-catalyzed hydration of alkenes.<sup>1</sup> Detailed interpretation of most of these reactions is limited, because the free energy change associated with the elementary steps is not known. Marcus pointed out that the natural comparison point for intrinsic reactivity for electron transfer (ET) reactions is the self-exchange-ET rate constant,  $k_{ii}$ .<sup>2</sup> He described the relationship between intrinsic reactivity and the rate constant for outer-sphere electron transfer for a net chemical reaction,  $k_{12}$ . For a reaction between a neutral species,  $i^0$ , and a different radical cation,  $j^+$ , eq 1,



electrostatic factors are unimportant, and the relationship is astonishingly simple, eq 2.<sup>2</sup>

$$k_{ij}(\text{calcd}) = (k_{ii}k_{jj}K_{ij}f_{ij})^{1/2} \quad (2a)$$

$$\ln(f_{ij}) = [\ln(K_{ij})^2 / [4 \ln(k_{ii}k_{jj}/Z^2)]] \quad (2b)$$

The self-exchange rate constants  $k_{ii}$  and  $k_{jj}$ , and the equilibrium constant,  $K_{ij}$ , are the principal parameters determining  $k_{ij}$ . A preexponential factor,  $Z$ , of  $10^{11} \text{ M}^{-1} \text{ s}^{-1}$  is often used; for our systems, the results are relatively insensitive to the value employed. The  $f_{ij}$  is 1 when  $K_{ij} = 1$  and decreases as the reaction becomes more exoenergetic; no reactions in our data set have  $f_{ij}$  less than 0.1. Formal potentials were determined using CV for all couples described here, and the  $K_{ij}$ 's used in eq 2 are derived from them. A more general form of eq 2 that includes the work terms necessary for multiply charged reactants, has been successfully applied to a wide variety of inorganic, organic, organometallic, and biochemical reactions.<sup>3–5</sup> However, studying couples having similar  $k_{ii}$  values principally tests the dependence of  $k_{ij}$  upon  $K_{ij}$ , which is known to work well.<sup>3–5</sup>

This work describes our results for a database of ET reactions, such as those described in eq 1, among a large set of compounds diverse in both intrinsic reactivity and structure, from which we have extracted reliable estimates of the intrinsic ET reactivity of all compounds in the set. The database allows an unprecedented test of structure–reactivity relationships and modern ET theory with some unexpected results, including close agreement with classical Marcus theory, striking steric effects, and estimates of the ET overlap integral and its dependence on reactant type and structure.

**Determination of  $\Delta G_{ii}^\ddagger$ (fit) Values.** We have measured  $k_{ij}$  using stopped-flow spectrophotometry for a diverse group of 0/1+ couples for which  $E^{0'}$  has been measured under the same conditions. Our principal contribution has been to include couples having a very large range of  $k_{ii}$  values. This has been achieved especially by including hydrazines of varied structural types, which causes a wide variation in the size of the geometry change between the oxidation states and, hence, the reorganization energy.<sup>6–11</sup> Having both large and small  $k_{ii}$  couples available for study greatly expands the range of accessible compounds, as compared to those that can be studied directly under self-exchange-ET conditions. The reaction is always at equilibrium under self-exchange-ET conditions, and electron exchange is usually detected using magnetic resonance (MR) line-broadening. However, the amount of broadening becomes too small to measure accurately for  $k_{ii}$  below  $\sim 7 \times 10^2 \text{ M}^{-1} \text{ s}^{-1}$ .<sup>12</sup> The upper limit for self-exchange-ET studies appears to be  $\sim 3 \times 10^9 \text{ M}^{-1} \text{ s}^{-1}$  (that for tetra-

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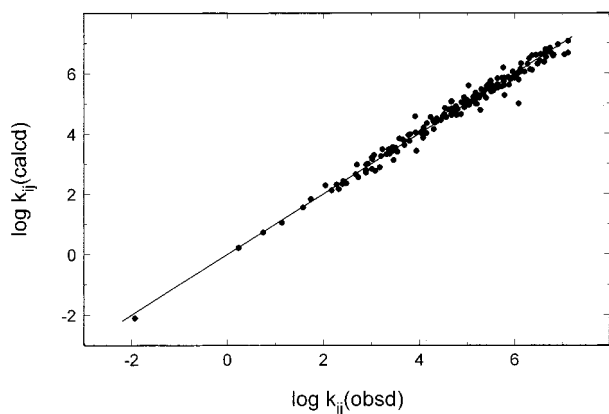


FIGURE 1. Plot of  $\log k_{ij}(\text{calcd})$  versus  $\log k_{ij}(\text{obsd})$  for the 164-reaction data set.

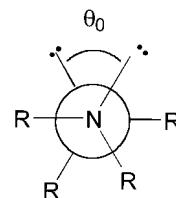
methyl-*p*-phenylenediamine, TMPD<sup>0/+</sup>).<sup>13</sup> Nearly two-thirds of the couples we have studied by stopped-flow (36 of 56) could not have their  $k_{ii}$  values measured under direct self-exchange conditions.

Stopped-flow kinetic data has frequently been used to establish relative reactivity by selecting a few isolable “oxidants” and determining the  $k_{ij}$  for their reactions with several reduced species. Extracting the  $k_{ii}$  for some compounds then required directly measuring  $k_{ii}$  for the others, typically using MR methods. In contrast, we do not use any directly measured self-exchange rate constants to extract  $k_{ii}$  from our  $k_{ij}$  data. Instead, the  $k_{ij}(\text{obsd})$  and  $E^0$  values, both measured at 25 °C in acetonitrile containing 0.1 M tetrabutylammonium perchlorate, are subjected to a least-squares fit of all of the data to eq 2 to extract  $k_{ii}(\text{fit})$  values.<sup>9–11</sup> The observed  $k_{ij}$  values are fit well using eq 2, as shown graphically in Figure 1. One hundred and fifty of the calculated  $k_{ij}$  values (92%) lie within a factor of 2 of the experimental value using the  $k_{ii}(\text{fit})$ s for the calculation, and an additional eight (5%) are within a factor of 3. Using eq 2 to obtain reliable  $k_{ii}(\text{fit})$ s requires measuring  $k_{ij}$  values between a wide range of the couples, which we have done.<sup>9–11</sup> The range of  $k_{ij}(\text{fit})$  values obtained is very large,  $2 \times 10^{14}$ , but we have found no dependence of the scatter of the  $k_{ij}(\text{fit})$ s on intrinsic reactivity,  $\Delta G^0$ , or type of reaction partner.

The  $k_{ii}(\text{fit})$  values have been converted to Eyring  $\Delta G_{ii}^{\ddagger}(\text{fit})$  values to allow their linear comparison, and we shall refer to these  $\Delta G_{ii}^{\ddagger}(\text{fit})$  values as intrinsic reactivity. For all of the reactions that we have been able to study, eq 2 works very well. Compounds of similar structure have similar intrinsic reactivity, and individual  $\Delta G_{ii}^{\ddagger}(\text{fit})$  established from a few reactions do not change much as more reactions are studied, suggesting that  $\Delta G_{ii}^{\ddagger}(\text{fit})$  is, indeed, determined to a reliability of a few tenths of a kcal/mol (the average  $\Delta\Delta G_{ii}^{\ddagger}(\text{fit}) = |0.592 \ln [k_{ij}(\text{obs})/k_{ij}(\text{fit})]|$  for the entire data set is 0.19 kcal/mol). Consequently, we believe they are useful for understanding the relationship between structure and reactivity. Moreover, there is a good agreement between  $\Delta G_{ii}^{\ddagger}(\text{fit})$  and directly measured self-exchange activation energy,  $\Delta G_{ii}^{\ddagger}(\text{self})$  for the 11 couples for which both have been determined, but the  $\Delta G_{ii}^{\ddagger}(\text{fit})$  values are systematically slightly larger than the  $\Delta G_{ii}^{\ddagger}(\text{self})$

values, and the average  $\Delta\Delta G_{ii}^{\ddagger} = \Delta G_{ii}^{\ddagger}(\text{fit}) - \Delta G_{ii}^{\ddagger}(\text{self}) = 0.6 \text{ kcal mol}^{-1}$ .<sup>9,10</sup> The  $\Delta G_{ii}^{\ddagger}(\text{fit})$  values arise from experimental rate constants for cross-reactions, so they must include the effect of averaging the preexponential factors for the cross-reactions studied, but  $\Delta G_{ii}^{\ddagger}(\text{self})$  values depend only on the preexponential factor for the self-exchange-reaction, and we believe that this may be the origin of this small difference. TMPD<sup>0/+</sup> is the lowest barrier and least sterically encumbered couple for which both values are available<sup>13b,c</sup> and has the largest  $\Delta\Delta G_{ii}^{\ddagger}$ , 1.5 kcal mol<sup>-1</sup>. The  $\Delta G_{ii}^{\ddagger}(\text{fit})$  value for TMPD<sup>0/+</sup> was determined from reactions with eight hydrazines: six tetra- $\alpha$ -branched, 22/*t*BuMe and 33NNMe<sub>2</sub>. None of the cross reactions is likely to have as good an orbital overlap between the reaction partners at the transition state as can TMPD<sup>0</sup> with TMPD<sup>+</sup>, and this might be the principal factor causing  $k_{ii}(\text{self})$  to be 13-fold larger than  $k_{ii}(\text{fit})$ . At the other reactivity extreme for which both barriers are available, the very hindered couple  $i\text{Pr}_2\text{N})_2^{0/+}$  has a  $\Delta\Delta G_{ii}^{\ddagger}$  within experimental error of zero.

**Factors Determining  $\Delta G_{ii}^{\ddagger}(\text{fit})$  Values.** Thirty-two hydrazine couples spanning a very large range of intrinsic reactivity, 16.5 kcal mol<sup>-1</sup> or 84% of the total  $\Delta G_{ii}^{\ddagger}(\text{fit})$  range observed, are listed in approximately descending order of their  $\Delta G_{ii}^{\ddagger}(\text{fit})$  in Table 1.<sup>14</sup> Abbreviations are used for the structures, and examples for the cyclic compounds are shown in Chart 1. A u indicates unsaturation  $\beta, \beta'$  to the hydrazine nitrogens, and k is used for a 3-oxo group in a 33N bicyclic ring. The ET reactivity of hydrazines falls into five groups on the basis of the connectivity of their substituents. The reason for this is that the lone-pair–

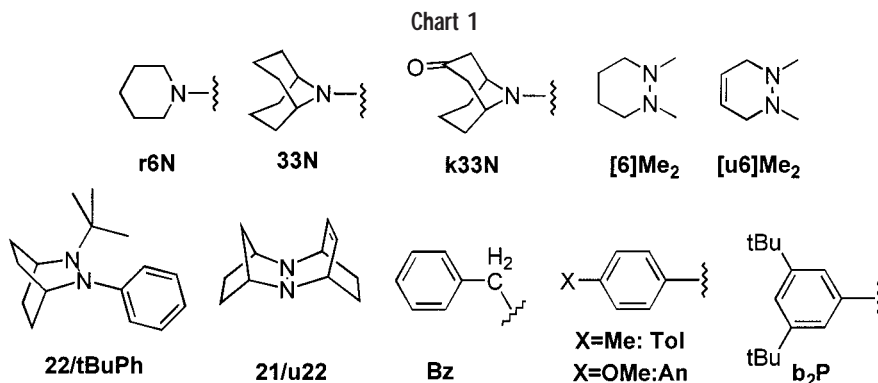


lone-pair twist angle about the NN bond of the neutral form,  $\theta_0$ , depends on the connectivity of the substituents.  $\theta_0$  is a primary structural feature influencing intrinsic reactivity, because it strongly affects the reorganization energy. Neutral hydrazines have significantly pyramidalized nitrogens<sup>14</sup> and electronically prefer  $\theta_0 \sim 90^\circ$  conformations, because this allows maximum lone pair,  $\sigma^*$  stabilizing interactions (often stated as avoiding overlap of the lone pair orbitals). The  $\theta_0$  values for the hydrazines were determined using photoelectron spectroscopic lone pair–lone pair orbital vertical ionization potential differences<sup>15</sup> and, in several cases, by X-ray crystallography.<sup>16</sup> Hydrazine radical cations have one antibonding ( $\pi^*$ ) and two bonding ( $\pi$ ) electrons and, therefore, have geometries that are very different from their related neutral molecules, showing a strong electronic preference for being untwisted ( $\theta_+$  near 0 or  $180^\circ$ ), and their nitrogens are considerably flattened (0 and  $180^\circ$  twists become the same when the nitrogens become planar).<sup>14</sup> The largest geometry changes and, hence, the largest reorganization energies and ET

**Table 1. Intrinsic Reactivity, NN Twist Angle Groups, Estimated  $\Delta G_{ii}^*$  Values, and Calculated Vertical Reorganization Enthalpies of Hydrazines<sup>a</sup>**

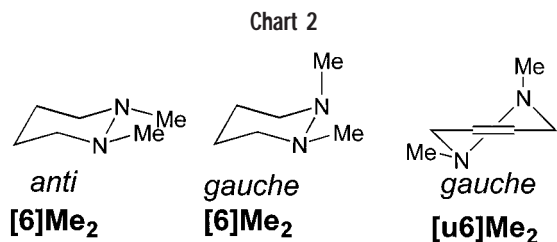
no.	couple	$\Delta G_{ii}^{\ddagger}(\text{fit})$	group ( $\theta_0$ )	$H_{ab}$ , $\Delta G_{ii}^*$ pairs	$\Delta H_v$
1	nPr <sub>2</sub> N) <sub>2</sub> <sup>0/+</sup>	21.9	1( $\theta_0 \sim 90^\circ$ )	<b>[0.01] 17.1</b>	
2	Et <sub>2</sub> N) <sub>2</sub> <sup>0/+</sup>	21.8	1	<b>[0.01] 17.0</b>	14.0 <sup>b</sup>
3	nHx <sub>2</sub> N) <sub>2</sub> <sup>0/+</sup>	21.4	1	<b>[0.01] 16.6</b>	
4	Bz <sub>2</sub> N) <sub>2</sub> <sup>0/+</sup>	21.1	1	<b>[0.01] 16.3</b>	13.5
5	iPrMeN) <sub>2</sub> <sup>0/+</sup>	20.2	1	0.026 – [16.6]	
6	iPr <sub>2</sub> NNMe <sub>2</sub> <sup>0/+</sup>	20.7	1	0.018 – [16.6]	
7	iPr <sub>2</sub> N) <sub>2</sub> <sup>0/+</sup>	21.0	1 <sup>c</sup>	<b>[0.01] 16.2</b>	13.9
8	cHx <sub>2</sub> N) <sub>2</sub> <sup>0/+</sup>	19.6	1 <sup>c</sup>	<b>[0.01] 14.8</b>	
9	nPr <sub>2</sub> NNMe <sub>2</sub> <sup>0/+</sup>	19.4	1	0.08 – [17.0]	
10	nPrMeN) <sub>2</sub> <sup>0/+</sup>	19.4	1	0.08 – [17.0]	
11	nBuMeN) <sub>2</sub> <sup>0/+</sup>	19.3	1	0.08 – [17.0]	
12	r7NNMe <sub>2</sub> <sup>0/+</sup>	18.2	1	0.21 – [17.0]	
13	Me <sub>2</sub> N) <sub>2</sub> <sup>0/+</sup>	17.3	1	0.45 – [17.0]	15.1
14	r6NNMe <sub>2</sub> <sup>0/+</sup>	16.8	1	[0.45]16.5	13.7
15	r5NNMe <sub>2</sub> <sup>0/+</sup>	16.7	1	[0.45]16.4	14.7
16	k33N) <sub>2</sub> <sup>0/+</sup>	15.2	2 <sup>c</sup> ( $\theta_0 = 180^\circ$ )	0.002 – [8.9]	
17	k33NN33 <sup>0/+</sup>	14.2	2	0.006 – [8.9]	
18	33N) <sub>2</sub> <sup>0/+</sup>	13.5	2 <sup>c</sup>	<b>[0.01] 8.9</b>	9.1 <sup>d</sup>
19	21/Me <sub>2</sub> <sup>0/+</sup>	16.1	3( $\theta_0 \sim 120^\circ$ )		
20	22/tBuPr <sup>0/+</sup>	15.8	3	<b>[0.001] 11.1</b>	10.5
21	22/tBuMe <sup>0/+</sup>	15.2	3		11.6
22	[u6]Me <sub>2</sub> <sup>0/+</sup>	17.4	only <i>gauche</i>		
23	[6]Me <sub>2</sub> <sup>0/+</sup>	15.0	<i>anti</i> > <i>gauche</i>		
24	22/22 <sup>0/+</sup>	14.7	4 <sup>c</sup> ( $\theta_0 \sim 0^\circ$ )	<b>[0.01] 10.0</b>	8.3 <sup>e</sup>
25	21/u22 <sup>0/+</sup>	13.4	4	<b>[0.01] 8.8</b>	9.7
26	22/u22 <sup>0/+</sup>	13.2	4 <sup>c</sup>	<b>[0.01] 8.6</b>	9.1
27	22/u23 <sup>0/+</sup>	12.8	4 <sup>c</sup>	<b>[0.01] 8.2</b>	9.6
28	21/21 <sup>0/+</sup>	12.7	4	<b>[0.01] 8.1</b>	8.8
29	22/tBuPh <sup>0/+</sup>	13.3	5 <sup>c</sup> ( <i>n</i> -aryl)		9.5
30	22/Ph <sub>2</sub> <sup>0/+</sup>	11.0	5		9.0
31	(b <sub>2</sub> P) <sub>2</sub> N) <sub>2</sub> <sup>0/+</sup>	8.2 <sup>f</sup>	5		
32	tol <sub>2</sub> N) <sub>2</sub> <sup>0/+</sup>	5.4	5		

<sup>a</sup>  $\Delta G_{ii}^{\ddagger}(\text{fit})$ ,  $\Delta G_{ii}^*$ , and  $\Delta H_v$  in kcal mol<sup>-1</sup>;  $\Delta H_v$  calculated using AM1. <sup>b</sup> This entry corresponds to using the minimum enthalpy neutral and cation conformations obtained. Values for various minimums of neutral and cation in other combinations range from 13.2 to 14.8 kcal mol<sup>-1</sup>. Obtaining such ranges is an increasing problem as conformational complexity increases. <sup>c</sup> Crystal structures available. <sup>d</sup> Constrained to  $\theta = 180^\circ$ ; a larger value is obtained using the AM1 optimum  $\theta$  values (which deviate from the experimental values of  $180^\circ$ ). <sup>e</sup> AM1 twist angle for the neutral is  $0^\circ$ , which is incorrect (it is really about  $15^\circ$ ), making this value too small relative to the other Group 4 entries. <sup>f</sup> Three of the reactions studied have  $k_{ii}(\text{fit})/k_{ii}(\text{obs})$  values lying outside the range 3 or 0.33 (so one-half of the six reactions studied that fit eq 2 poorest involve this couple). The structure of this compound was given incorrectly in ref 11.



barriers ( $\Delta G_{ii}^{\ddagger}(\text{fit}) \geq 16.7$  kcal mol<sup>-1</sup>) occur for Group 1,  $\theta_0 = \sim 90^\circ$ , hydrazines, entries 1–15 (e.1–15) of Table 1. Group 2,  $\theta_0 = 180^\circ$ , compounds (e.16–18) result when the substituents at nitrogen become large enough and also are unable to rotate away from N–C<sub>α</sub> twist angles that give large nonbonded steric interactions between α-substituents. Just having four α-branched substituents, as for iPr<sub>2</sub>N)<sub>2</sub> (e.7) and c-Hx<sub>2</sub>N)<sub>2</sub> (e.8), is not sufficient to make  $\theta_0 = 180^\circ$  conformations predominate, because their alkyl groups can rotate out of conformations having large steric interactions with substituents on the other nitrogen. The Group 3 hydrazines (e.19–21) are forced to  $\theta_0 \sim 120^\circ$ ,

because they have near- $0^\circ$  bicyclic ring CNNC angles imposed by their structures. The electronically least favorable  $\theta_0 \sim 0^\circ$  conformations are formally possible if the alkyl groups are syn to each other, but a syn conformation has never been observed in a bicyclic hydrazine. In contrast, the chair ring of dimethylhexahydropyridazine, [6]Me<sub>2</sub>, allows both anti ( $\theta_0 \sim 180^\circ$ ) and gauche ( $\theta_0 \sim 60^\circ$ ) conformations to be occupied (Chart 2); they differ in energy by only 0.2 kcal mol<sup>-1</sup>.<sup>17</sup> The intrinsic reactivity for [6]Me<sub>2</sub><sup>0/+</sup> is slightly less than that of the  $\theta_0 \sim 180^\circ$  33N)<sub>2</sub><sup>0/+</sup>, which also has flatter nitrogens because of its α-branched carbons. Despite differing in structure

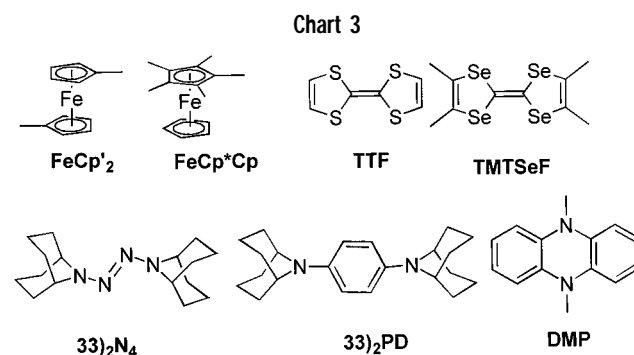


by only two hydrogens, dimethyltetrahydro-pyridazine  $[\text{u6}]\text{Me}_2$  can be detected in only the gauche conformation<sup>17</sup> and we argue, as a result, has reactivity close to that of the unhindered Group 1 hydrazine  $\text{Me}_2\text{N}_2$ . Group 4 hydrazines have both CNNC angles forced to be near  $0^\circ$ , which produces syn lone pair neutral conformations having  $\theta_0 \sim 0^\circ$ , and nonbonded steric interactions flatten their nitrogens relative to the  $\theta_0 = 180^\circ$  Group 2 ones. Group 4 hydrazines, therefore, have especially small geometry changes upon electron loss and are the lowest  $\Delta G_{ij}^{\ddagger}(\text{fit})$  hydrazines that lack aryl substitution, and their ET could be studied by direct self-exchange methods as well as stopped-flow.<sup>12,18</sup> The highest barrier couple of this group is 22/22, shown by crystallography to be twisted  $15^\circ$  in the neutral form.<sup>16b</sup> A twisted neutral structure increases the vertical reorganization energy,  $\lambda_v$ , because twisting the cation radical is difficult. The other saturated compounds are untwisted, in agreement with calculated geometries, and have their  $\Delta G_{ij}^{\ddagger}(\text{fit})$  values in the order of their calculated  $\lambda_v$  values for electron loss.<sup>18</sup> Decreasing ring size makes the nitrogens more pyramidal in both oxidation states, and the nitrogens are pyramidal enough for 21/21<sup>+</sup> that the barrier to their becoming planar is 4.6 kcal mol<sup>-1</sup>, measured by ESR.<sup>19</sup> Successively replacing alkyl substituents by aryl ones (Group 5) substantially lowers the ET barrier. For example, replacing the isopropyl of 22/tBu-iPr (e.20) by phenyl to give 22/tBuPh, (e.29) lowers  $\Delta G_{ij}^{\ddagger}(\text{fit})$  by 2.5 kcal mol<sup>-1</sup>, and a second aryl substitution in going to 22/Ph<sub>2</sub> (e.30) lowers  $\Delta G_{ij}^{\ddagger}(\text{fit})$  an additional 2.3 kcal mol<sup>-1</sup>. The tetra-arylhydrazines, e.31 and e.32, are the most reactive hydrazines studied. The  $\Delta G_{ij}^{\ddagger}(\text{fit})$  for  $\text{tol}_2\text{N}_2$  (e.32) lies between those for the unhindered nearly planar aminoaromatic couples, TMPD and DMP. Turning to nonhydrazine couples in Table 2 and Chart 3, the bicyclononyl protected 2-tetrazenes,  $(33)_2\text{N}_4$  (e.33) and  $(33)_2\text{-PD}$  (e.34), are intermediate in reactivity between hydrazines and aromatic compounds. They are much more reactive than their hydrazine counterparts:  $\Delta \Delta G_{ij}^{\ddagger}(\text{fit}) = 3.7$  kcal mol<sup>-1</sup> for the 33N-substituted systems (e.18 and e.34) and 4.5 kcal mol<sup>-1</sup> for the keto systems (e.16 and e.33). The four ferrocenes, e.35–38, lie in a narrow reactivity range and are less reactive than any of the aromatic organic compounds. Aromatic compounds make up the upper quartile of ET reactivity measured. The  $k_{ii}(\text{fit})$  values for the fastest ones are near or at the diffusion limit: TMTSF<sup>0/+</sup>,  $1.2 \times 10^{11}$ ; An<sub>3</sub>N<sup>0/+</sup>,  $2.4 \times 10^{10}$ ; and TTF<sup>0/+</sup>,  $1.4 \times 10^{10}$  M<sup>-1</sup> s<sup>-1</sup> (e.46–44); so their intrinsic reactivity can be accurately determined only by using cross-reactions. The  $k_{ij}(\text{fit})$  obtained for An<sub>4</sub>PD<sup>0/+</sup> (e.42),  $3.6 \times 10^8$  M<sup>-1</sup> s<sup>-1</sup>, is more than a factor of 100 less than

**Table 2. Intrinsic Reactivity for Non-Hydrazine Couples<sup>a</sup>**

no.	couple	$\Delta G_{ij}^{\ddagger}$	$\Delta G_{ij}^{*b}$	$(\Delta H_v + 2)^c$	type
33	$\text{k33})_2\text{N}_4^{0/+}$	10.7			2-tetrazene
34	$\text{33})_2\text{N}_4^{0/+}$	9.8			2-tetrazene
35	$\text{FeCp}^*_2^{0/+}$	8.3			ferrocene
36	$\text{FeCp}^*\text{Cp}^{0/+}$	8.0			ferrocene
37	$\text{FeCp}^{*2}^{0/+}$	7.9			ferrocene
38	$\text{FeCp}_2^{0/+}$	7.7			ferrocene
39	$\text{k33})_2\text{PD}^{0/+}$	7.2			PD
40	TMPD <sup>0/+</sup>	6.5	6.58	6.67	PD
41	$\text{33})_2\text{PD}^{0/+}$	6.2			PD
42	An <sub>4</sub> PD <sup>0/+</sup>	5.8			triarylamine/PD
43	DMP <sup>0/+</sup>	5.3	5.43	5.88	aromatic
44	TTF <sup>0/+</sup>	3.6	3.83	3.73	aromatic
45	An <sub>3</sub> N <sup>0/+</sup>	3.3			triarylamine
46	TMTSF <sup>0/+</sup>	2.3	2.6 <sup>d</sup>		aromatic

<sup>a</sup>  $\Delta G_{ij}^{\ddagger}(\text{fit})$ ,  $\Delta G_{ij}^{*b}$  and  $\Delta H_v$  in kcal mol<sup>-1</sup>. <sup>b</sup> Calculated with  $H_{ab} = 0.5$ ,  $K_e = 1$  using eqs 3–5. <sup>c</sup> AM1 calculated, plus 2 kcal mol<sup>-1</sup> as an estimate of  $\lambda_s/4$ . <sup>d</sup> If a 2 kcal mol<sup>-1</sup> solvent contribution to  $\Delta G_{ij}^{*b}$  is assumed, this leaves only 0.6 kcal mol<sup>-1</sup> for  $\Delta H_v$ .



that for An<sub>3</sub>N<sup>0/+</sup> but close to that obtained under self-exchange conditions by NMR in CDCl<sub>3</sub>,  $3.5(3) \times 10^8$ .<sup>22</sup>

**Electronic Coupling and Reorganization Energy Effects on Intrinsic Reactivity.** It is revealing to consider the observed intrinsic reactivity using more modern ET theory. The fundamental assumption giving rise to eq 2 is that the reactions are activation-barrier-limited, so the barrier for the cross-reaction will be the average of those for the related self-exchange reactions.<sup>2</sup> Marcus obtained eq 2 by assuming that ET reactions are adiabatic, and preexponential factors should be nearly constant for adiabatic reactions. The adiabatic preexponential factor is often approximated as  $3 \times 10^{10} h\nu_v$ , where  $h\nu_v$  is the energy corresponding to the inherent barrier-crossing frequency, usually thought of as a bond-stretching mode. The  $h\nu_v$  values for all of the couples studied here are believed to vary between  $\sim 400$  and  $1600$  cm<sup>-1</sup>. However, achieving adiabaticity requires rather large  $H_{ab}$  values, which are not believed to occur for most intermolecular ET reactions, and certainly not for ones between compounds as hindered as for most of the cross-reactions reported here. In a recent review on ET, Bixon and Jortner said that although there was “lively discussion” in the 1960s about whether ET reactions were adiabatic or nonadiabatic, it has now been established that the great majority are nonadiabatic.<sup>21</sup> But if this were the case for our reactions, eq 2 would not work well, because nonadiabatic reaction rate constants are not only controlled by activation barriers (or Franck–Condon factors using



Jortner's vibronic coupling theory) but also by widely varying preexponential factors. The nonadiabatic preexponential factor is proportional to the ET orbital overlap integral squared,  $(H_{ab})^2$  and is also directly proportional to  $e^{-S}$ , where  $S$  is the vibronic coupling constant (Huang–Rhys factor), which is the ratio of  $h\nu_v$  to the internal vibrational component of  $\lambda$ ,  $\lambda_v$ . Tetraalkylhydrazines have unusually large  $S$  values, exceeding 20, but aromatic compounds and ferrocenes have much smaller  $S$  values, certainly under 6, so simply from the  $e^{-S}$  term, a factor of above a million-fold faster ET should occur for aromatic compounds and ferrocenes than for hydrazines. The consequences of such an effect on the preexponential term are not observed experimentally. The same  $k_{ii}$  value suffices to calculate  $k_{ij}$  whether a hydrazine couple is reacted with a ferrocene, an aromatic compound, or another hydrazine. Equation 2 would not fit our data so well if the preexponential factor were sensitive to  $S$ .

To interpret our data, we modify the simplest non-adiabatic rate equation of Levich and Dogonadze.<sup>21</sup> The modifications are (a) the ET activation free energy,  $\Delta G^*$ , of classical Marcus–Hush two-state theory (eq 3)<sup>3</sup> replaces  $\lambda/4$  to allow reactions that do not have vanishingly small  $H_{ab}$  to be treated

$$\Delta G^* = \lambda/4 - H_{ab} + (H_{ab})^2/\lambda \quad (3)$$

and (b) the encounter complex formation constant,  $K_e$ , has been inserted so that the equation can be used for intermolecular ET. This results in eq 4, which like adiabatic Marcus theory, uses only  $\lambda$  and  $H_{ab}$  to predict the rate constant.

$$k_{L\&D}(25\text{ }^\circ\text{C}) = (1.52 \times 10^{14})(K_e H_{ab}^2/\lambda^{1/2}) \exp[-\Delta G^*/RT] \quad (4)$$

Importantly, it is necessary to know both  $H_{ab}$  and  $K_e$  to extract  $\Delta G^*$  (and  $\lambda$ ) from intrinsic reactivity. There is, however, no way of experimentally determining  $K_e$  for our reactions, and factors that raise  $H_{ab}$  appear likely to raise  $K_e$ , as well. Interpretation of intermolecular reactions must include both factors, so we shall use  $H'_{ab} = K_e^{1/2} H_{ab}$  (units,  $M^{-1/2}$  kcal mol<sup>-1</sup>) in discussing our results.<sup>11</sup> Substitution of  $k_{ij}(\text{fit})$  for  $k_{L\&D}(25\text{ }^\circ\text{C})$  in eq 4 leads to eq 5.

$$\Delta G^*_{ii} = 0.592[32.655 - \ln(k_{ij}(\text{fit})) + \ln(H'_{ab})^2/\lambda^{1/2}] \quad (5)$$

We believe that eqs 3–5 provide an internally consistent interpretation of our data and allow important conclusions to be drawn about the reactions studied.

Figure 2 shows an example of the dependence of  $\lambda/4$  and  $\Delta G^*_{ii}$  on  $H'_{ab}$ . The shape of the curve is similar to that obtained using an adiabatic rate equation with an electronic transmission coefficient in the preexponential term, but using eq 4 requires fewer parameters, because one does not need to know the  $h\nu_v$  and  $\lambda_v$  to calculate the preexponential term.<sup>11</sup> All of the points on the solid line of Figure 2 correspond to the same  $\Delta G^*_{ii}(\text{fit})$  value, emphasizing that if intrinsic reactivity correlates with  $\lambda$ ,  $H'_{ab}$  cannot be changing very much. The curved lines of

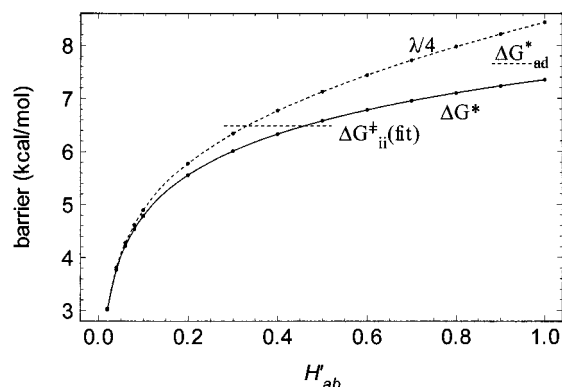


FIGURE 2. Plot of  $\Delta G^*$  (using eqs 3–5) and the corresponding  $\lambda/4$  (using eq 3 with  $H_{ab}$  replaced by  $H'_{ab}$ ) for the  $k_{ij}(\text{fit})$  of  $\text{TMPD}^{0/+}$  ( $\Delta G^*_{ii}(\text{fit}) = 6.5$  kcal/mol). An  $h\nu_v$  of  $1500\text{ cm}^{-1}$  was used to calculate  $\Delta G^*_{ad}$ .

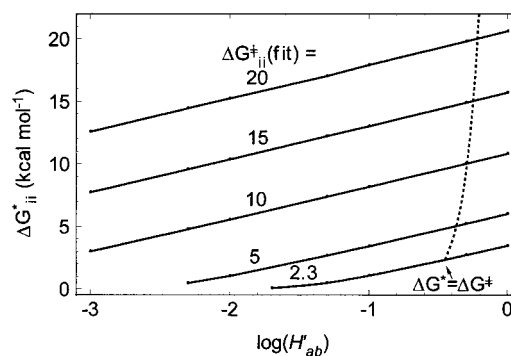


FIGURE 3. Plots of  $\Delta G^*_{ii}$  versus  $\log(H'_{ab})$  for various values of  $\Delta G^*_{ii}(\text{fit})$ .

Figure 2 become nearly linear when  $\log(H'_{ab})$  is used as the x axis, as shown in Figure 3, and the lines for various  $\Delta G^*_{ii}(\text{fit})$  values are displaced by approximately the difference in  $\Delta G^*_{ii}(\text{fit})$ . The broken line in Figure 3 shows  $\log(H'_{ab})$  values for which  $\Delta G^*_{ii}$  is equal to  $\Delta G^*_{ii}(\text{fit})$  according to eq 5, dropping from about  $H_{ab} = 0.63$  for the largest  $\Delta G^*_{ii}(\text{fit})$  ( $\sim 22$  kcal mol<sup>-1</sup> for nR<sub>4</sub>N<sub>2</sub>) to 0.35 for the smallest (2.3 kcal mol<sup>-1</sup> for TMTSF). It may be seen that  $\Delta G^*_{ii}$  is significantly less than  $\Delta G^*_{ii}(\text{fit})$  for high-barrier compounds that have small  $H_{ab}$ , but for less-hindered compounds that have higher  $H_{ab}$  values, mostly the aromatics,  $\Delta G^*_{ii}$  will approach  $\Delta G^*_{ii}(\text{fit})$ .

Which  $H_{ab}$  values are appropriate for intermolecular ET reactions has not been clear. For what is probably the most discussed example, Grampp and Jaenicke suggest that the  $\text{TMPD}^{0/+}$  transition state has a geometry with the aryl rings in  $\pi$ -stacking contact with parallel long axes, as they are in crystals, and that a rather low value of  $H_{ab} = 0.1$  kcal mol<sup>-1</sup> is necessary to be consistent with the directly measured  $k_{ij}$  value,<sup>13</sup> but Rauhut and Clark have calculated values using ab initio theory for transition states having a wider range of geometries, estimating an  $H_{ab}$  of  $0.65$  kcal mol<sup>-1</sup>.<sup>22</sup> Weaver and co-workers have argued from solvent studies on self-exchange-ET rate constants of metallocenes that ferrocene and decamethylferrocene have  $H_{ab}$  values of  $0.1$  and  $0.2$  kcal mol<sup>-1</sup>, respectively.<sup>23</sup> We consider our data on these and related molecules using the  $\Delta G^*_{ii}$  versus  $\log(H'_{ab})$  plot of Figure 4. These are

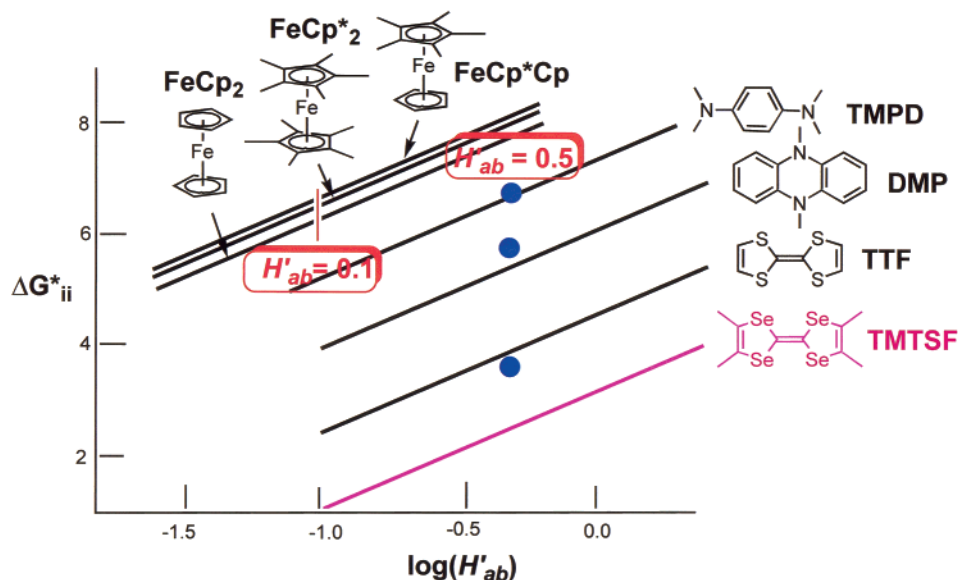


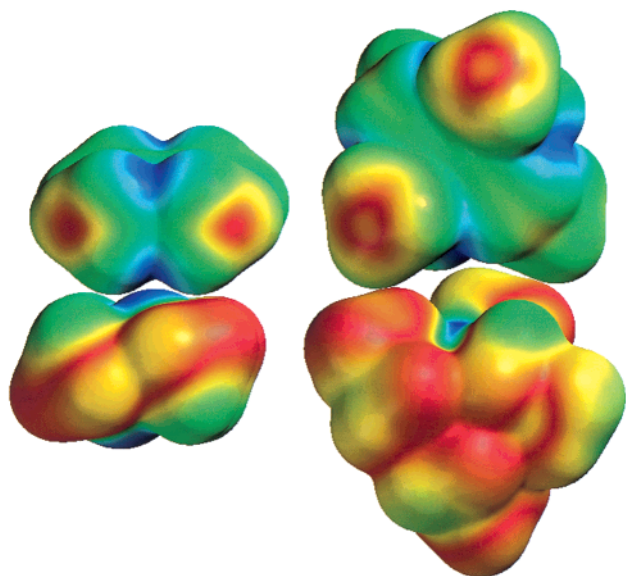
FIGURE 4. Plots of  $\Delta G_{ii}^{\ddagger}$  versus  $\log(H'_{ab})$  for some ferrocene and aromatic couples.

the least-hindered compounds studied, so their  $H_{ab}$  values should be the highest. The enthalpy portion of  $\lambda_v$  can be calculated using the AM1 semiempirical method, and rather good correlation between the calculated contribution to the barrier (one-quarter of the calculated internal reorganization enthalpy), which we will call  $\Delta H_v$ , and  $\Delta G_{ii}^{\ddagger}(\text{fit})$  has been noted.<sup>10,11</sup> In Figure 4, we highlight an  $H_{ab}$  value of 0.5 (x axis value of  $-0.3$ ) for the planar aromatics, where the  $\Delta H_v$  values for  $\text{TMPD}^{0/+}$ ,  $\text{DMP}^{0/+}$ , and  $\text{TTF}^{0/+}$  plus 2 kcal mol<sup>-1</sup> (equivalent to using  $\lambda_s = 8$  kcal mol<sup>-1</sup> for these compounds) lie close to the experimental data lines. If the  $H_{ab}$  values for these couples differed by much more than a factor of 2 (an increment of 0.3 in the  $\log(H_{ab})$  x axis value), one would expect the separations of the  $\Delta G_{ii}^{\ddagger}$  values to become significantly different from those of the  $\Delta H_v$  values. Using  $H_{ab} = 0.5$  is also close to using Clark's ab initio calculated  $H_{ab}$  value for  $\text{TMPD}^{0/+}$  ( $K_e$  is estimated to be somewhat less than 1 M<sup>-1</sup>).<sup>13</sup> Another reason for choosing  $H_{ab} = 0.5$  as an estimate for these compounds is that the fastest couple studied,  $\text{TMTSF}^{0/+}$  (for which  $\Delta H_v$  is not available using AM1), has a  $\Delta G_{ii}^{\ddagger}(\text{fit})$  of only 2.3 kcal mol<sup>-1</sup>, and using  $H_{ab} = 0.5$  produces a  $\Delta G_{ii}^{\ddagger}$  of 2.6 kcal mol<sup>-1</sup>. If, as before, the solvent reorganization contribution is estimated to be 2 kcal mol<sup>-1</sup>, only 0.6 kcal mol<sup>-1</sup> remains for bond reorganization. Consequently, a smaller value of  $H_{ab}$  would not be reasonable. Figure 4 also demonstrates that the  $\Delta G_{ii}^{\ddagger}(\text{fit})$  values for ferrocenes require smaller  $H_{ab}$  values than those for aromatic compounds. If  $H_{ab} = 0.1$  is used for the ferrocenes, as highlighted, their  $\lambda_v$  values are on the order of 16 kcal mol<sup>-1</sup> (using  $\lambda_s$  of 8 kcal/mol again) which is significantly larger than  $\lambda_v$  has been argued to be,<sup>23</sup> so an even smaller  $H_{ab}$  value might be appropriate.

In contrast to planar aromatics, the  $\Delta G_{ii}^{\ddagger}(\text{fit})$  values for Group 1 hydrazines imply a substantial range of  $H_{ab}$  values. Double *n*-alkyl replacement of methyl on  $\text{Me}_2\text{N}^{0/+}$ , going to  $\text{nPrMeN}^{0/+}$ ,  $\text{nPr}_2\text{NNMe}_2^{0/+}$ , and  $\text{nBuMeN}^{0/+}$  raises  $\Delta G_{ii}^{\ddagger}(\text{fit})$  by 2.0–2.1 kcal mol<sup>-1</sup>, and replacing all

four methyls in  $\text{Et}_2\text{N}^{0/+}$ ,  $\text{nPr}_2\text{N}^{0/+}$  and  $\text{nHx}_2\text{N}^{0/+}$  raises it by 4.1–4.6 kcal mol<sup>-1</sup>.<sup>24</sup> However,  $\lambda_v$  for *n*-alkyl compounds should be no larger than for  $\text{Me}_2\text{N}^{0/+}$  (we note that AM1 calculations get  $\Delta H_v$  to be slightly lower when groups larger than methyl are present; see the last column of Table 1), so their  $\Delta G_{ii}^{\ddagger}$  values should be the same or slightly lower than that of  $\text{Me}_2\text{N}^{0/+}$ . Moreover, a slightly smaller  $\lambda_s$  is expected as the alkyl group size increases, yet  $\text{Me}_2\text{N}^{0/+}$  is more than >2000 times more reactive than  $\text{Et}_2\text{N}^{0/+}$ . We believe that  $H_{ab}$  may be as large for  $\text{Me}_2\text{N}^{0/+}$  as it is for aromatics (ca. 0.5) because of the direct HOMO–LUMO overlap possible, as depicted in Figure 5.

This, in turn, requires that  $H_{ab}$  drop to  $\sim 0.1$  for the di-*n*-alkyl compounds and 0.01 for  $\text{nR}_2\text{N}^{0/+}$  to account for their lower reactivity in the absence of increased  $\Delta H_v$ . We suggested that these substantial changes in  $H_{ab}$  arise from a steric effect.<sup>10,11</sup> Nonbonded interactions between *n*-alkyl groups make the methyl groups of  $\text{Et}_2\text{N}^{0/+}$  assume alternating positions above and below the nitrogens and  $\alpha$ -carbons, which lie nearly in a plane (see Figure 5). This geometry precludes direct overlap of the high-density region of the LUMO near nitrogen of  $\text{Et}_2\text{N}^{0/+}$  with the highest density region of the HOMO of  $\text{Et}_2\text{N}^0$ . In striking contrast, the LUMO of  $\text{Me}_2\text{N}^{0/+}$  can achieve significant direct HOMO overlap near nitrogen with an approaching neutral molecule, resulting in a significantly larger  $H_{ab}$ . Our data show that no significant increase in  $\Delta G_{ii}^{\ddagger}(\text{fit})$  occurs for further enlarging the alkyl groups either by lengthening the alkyl chain or by  $\alpha$ -branching. For example,  $\text{iPr}_2\text{N}^{0/+}$  has a 0.8 kcal mol<sup>-1</sup> lower value than does  $\text{Et}_2\text{N}^{0/+}$ , which may reflect a lower  $\lambda_v$ .<sup>9</sup> If “leveling” to a nearly constant  $H_{ab}$  did not occur for larger alkyl groups,  $\Delta G_{ii}^{\ddagger}(\text{fit})$ s would be unlikely to track calculated  $\lambda_v$  values as they do, since the difference between  $\Delta G_{ii}^{\ddagger}(\text{fit})$  and  $\lambda/4$  gets larger as  $H_{ab}$  decreases, as shown in Figure 2. We examine the postulate that  $H_{ab}$  becomes nearly constant for unlinked *n*-alkyl and branched alkyl groups



**FIGURE 5.** HOMO electron density function mapped on the surface density function for Me<sub>2</sub>N)<sub>2</sub><sup>0</sup> (lower left) and the LUMO mapped on the surface density function of Me<sub>2</sub>N)<sub>2</sub><sup>+</sup> (upper left). Blue indicates the region of highest density for the mapped function; red is zero density. The equivalent functions are displayed for the Et<sub>2</sub>N)<sub>2</sub><sup>0</sup>–Et<sub>2</sub>N)<sub>2</sub><sup>+</sup> reactant pair on the right. Geometry optimization and electron density function maps were computed at the DFT B3LYP/6-311G\* level using Spartan 02.

in Table 1 by showing the  $\Delta G_{ii}^*$  value calculated by fixing  $H_{ab} = 0.01$  for all 13 such compounds. This results in the estimated  $\Delta G_{ii}^*$  values shown in bold, and a relatively constant  $\Delta G_{ii}^*$  of 16–17 kcal mol<sup>-1</sup> arises for the similar compounds (e.1–4; 7). For Group 1 compounds that do not have four *n*-alkyl groups (i.e., they have some methyl substituents), we calculated  $H_{ab}$  values from an assumed  $\Delta G_{ii}^*$  of 17.0 kcal/mol for *n*-alkyl groups and slightly less for branched alkyl groups and connected the numbers by reverse arrows (←), indicating that  $H_{ab}$  is calculated from the assumed  $\Delta G_{ii}^*$ . This produces intermediate  $H_{ab}$  estimates for intermediate numbers of *n*-alkyl groups in agreement with the anticipated steric effects. The N,N-ring compounds r5NNMe<sub>2</sub><sup>0/+</sup> and r6NNMe<sub>2</sub><sup>0/+</sup> have *n*-alkyl substituents but cannot attain the  $\pi$ -system blocking conformations shown in Figure 5, because the ring restricts alkyl group motion and  $H_{ab}$  of 0.45, like Me<sub>2</sub>N)<sub>2</sub><sup>0/+</sup>, is assumed; the more flexible ring system in r7NNMe<sub>2</sub><sup>0/+</sup> allows limited alkyl group blocking. This approach produces  $H_{ab}$ 's consistent with calculated  $\lambda_v$ 's and that drop substantially when HOMO–LUMO overlap near nitrogen is sterically precluded and ET must proceed through lower-density overlap near alkyl groups. However, it does not appear to drop below about 0.01 for Group 1 hydrazines. The vertical reorganization enthalpies for ET calculated by AM1, shown as  $\Delta H_v$ , are in remarkably good agreement with the experimental barriers using the  $H_{ab} = 0.01$  approximation for Group 3 and 4 hydrazines, too, especially considering how crude these calculations are. Group 2 compounds are expected to have very similar  $\lambda_v$ 's, and a decrease in  $H_{ab}$  by a factor of  $\sim 2$  per keto group is necessary to account for the observed increases in  $\Delta G_{ii}^*$

(fit) within this group relative to 33N)<sub>2</sub><sup>0/+</sup>. This is consistent with a hypothesis that contact for regions near the electron-withdrawing carbonyl groups is ineffective for producing  $H_{ab}$  because of low spin density in the radical cations.<sup>9–11</sup> If the two tetraarylhya-zines studied have about the same  $\Delta G_{ii}^*$ , the 2.8 kcal mol<sup>-1</sup> difference in their  $\Delta G_{ii}^*$ (fit) values would correspond to a 12-fold smaller  $H_{ab}$  value for the hindered (b<sub>2</sub>P)<sub>2</sub>N)<sub>2</sub><sup>0/+</sup>, which is consistent with an increase in  $H_{ab}$  for replacement of branched alkyl by an unhindered aryl group.

**Comparisons with Gas-Phase Reactions.** Equation 2 also has been successfully applied to 40 gas-phase ET reactions among some of the hydrazines in acetonitrile that were just described, demonstrating not only that these reactions are activation-limited, but also that solvent reorganization energy is not especially important when vibrational reorganization energy is large.<sup>25</sup> The same approach to extracting intrinsic reactivity for gas-phase reactions was used with similar internal agreement. Reactivity was similar to that found in solution with an exception for the reversal of reactivity of [u6]Me<sub>2</sub> and [6]-Me<sub>2</sub>, which was attributed to the much higher cation–neutral association energy in the gas-phase precursor complex.

## Conclusion

The classical Marcus cross-rate expression, eq 2, correlates all of our data extraordinarily well, implying that these reactions are activation-barrier-limited. These experiments allow accurate estimation of intrinsic reactivity for a wide array of compounds ranging in structure from TMTSF to nPr<sub>2</sub>N)<sub>2</sub> and spanning a  $k_{ij}$  range of  $2 \times 10^{14}$ . The data set includes many reactions between very hindered compounds that appear limited to transferring an electron via nonbonded contact of saturated alkyl groups. The correlation observed demonstrates that the same intrinsic rate constants suffice whether a couple is reacted with a hydrazine, a ferrocene, or an aromatic compound, and that preexponential factors as well as the barriers for the cross reactions must effectively average. Such behavior is not expected from the perspective of nonadiabatic vibronic coupling ET theory.<sup>21</sup> Thus, although the effective electronic coupling for electron transfer,  $H_{ab}$ , appears to be on the order of only 0.01 for many of the couples studied, their electron transfer cannot be considered “nonadiabatic” in the sense the term is currently used. Although we are unable to make an unequivocal determination of  $H_{ab}$ , a combination of rate data, computational results, and available structural data has allowed us to place relatively narrow boundaries on what  $H_{ab}$  may be for a wide array of compounds.

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